

The header features a blue and orange horizontal bar. Above it, there are three panels: a green and blue grid-like pattern on the left, a molecular structure with red and yellow spheres in the center, and a vertical grey and white striped pattern on the right.


Fast ion ignition with ultraintense laser pulses

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The footer consists of a blue and orange horizontal bar at the bottom of the slide.



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Motivation

Fast Ignition with energetic ions

- TNSA scheme with thin foils
- Ion acceleration in underdense plasmas, collisionless shocks
- RPA scheme in the hole-boring mode
 - Analytical estimates
 - 1D/ 2D PIC simulation results
 - Fast ignition requirements
- Cone-guided carbon acceleration
- Large-scale fusion pellet with *in-situ* DT acceleration

Conclusion

Fast Ignition in Inertial Fusion Energy production



M. Tabak et al., Phys. Plasmas 1994

M. Tabak et al., Fusion Science and Technology 2006

fuel gain in an isobaric configuration:

$$G = \phi Q M / (\eta E_{driver}), \quad \phi = \frac{h}{h + 7g \text{cm}^{-2}}$$

$$\eta E_{driver} = E_{hs} + E_{M fuel}$$

$$E_{hs} [\text{MJ}] = M_{hs} C_v T_{hs} = \frac{4\pi}{3} \frac{h_{hs}^3}{\rho_{hs}^2} C_v T_{hs}$$

$$E_{M fuel} [\text{MJ}] = 0.3\alpha (M - M_{hs}) \rho_{M fuel}^{2/3}$$

Ignition requirements for FI

S. Atzeni, Phys. Plasmas 1999

- Laser propagation in overcritical plasma: hole boring, relativ. transparency

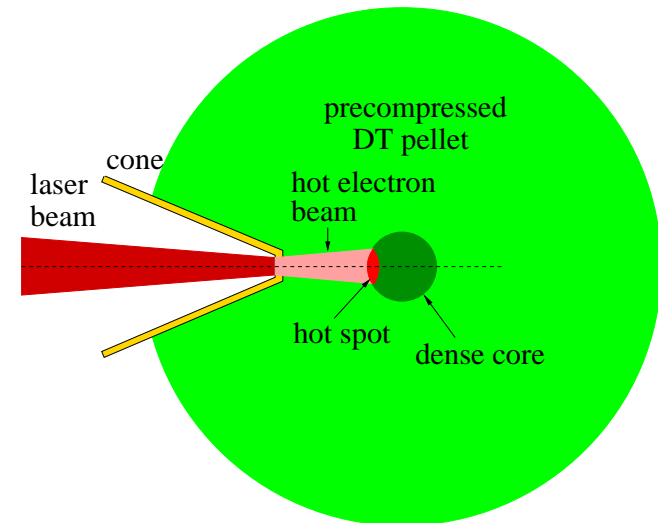
high laser intensities to propagate over mm
plasma length in 10 - 20 ps

- Cone-guided electron acceleration

R. Kodama et al., Nature 2002

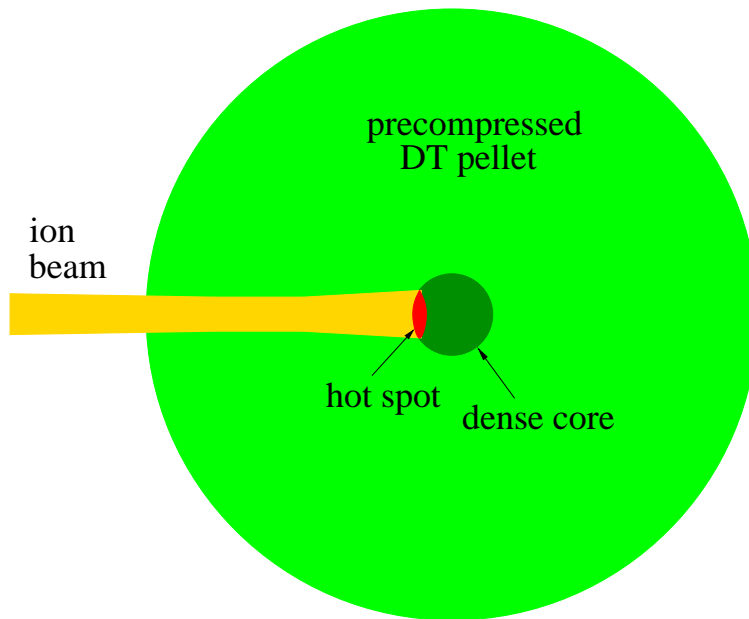
A. Debayle et al., J. Phys. Conf. Ser. 2010

higher coupling efficiency to e^- ,
problems in transport to the compressed fuel
influence of the implosion process



Fast Ignition by laser-accelerated proton/ light ion beams

M. Roth et al., PRL 2001
V.Yu. Bychenkov et al., Plasma Phys. Rep. 2001
S. Atzeni et al., Nucl. Fusion 2002
M. Temporal et al., Phys. Plasmas 2002
J.F. Fernandes et al., Nucl. Fusion 2009



less deflection by atomic scatters, induc. fields
divergence < 10 mrad required
spread in energy → spread in arrival times

$$E_{ig} = 90(d_{\text{mm}})^{0.7} \left(\frac{\rho}{100 \text{g cm}^{-2}} \right)^{-1.3} \text{ kJ}$$

$$T_p = 1 - 8 \text{ MeV}$$

NOVA-Petawatt:

$$I_L = 5 \times 10^{20} \text{ W cm}^{-2}, 60 \text{ MeV protons}$$

$$\eta = 10 - 20\%, \text{ emittance} < 0.004 \text{ mm mrad}$$

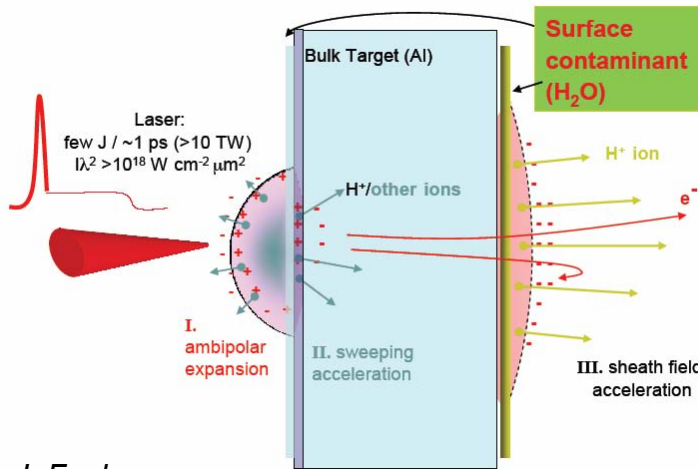
focused beams by shaped targets
focusability at high currents ?

$$\varepsilon_p = 10 \text{ MeV}, r_{hs} \sim 15 \mu\text{m}, \tau_{dep} \sim 10 \text{ ps}$$

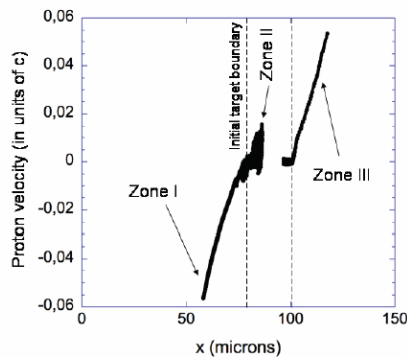
$$E_{pb} = 50 \text{ kJ} \rightarrow 10^{22} \text{ protons/cm}^2$$

TNSA of protons and light ions (foil targets)

Interaction scheme



J. Fuchs,
 HDR thesis 2010



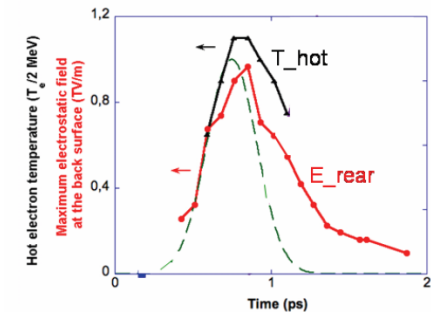
isothermal expansion

Gurevich 1966, Mora 2003

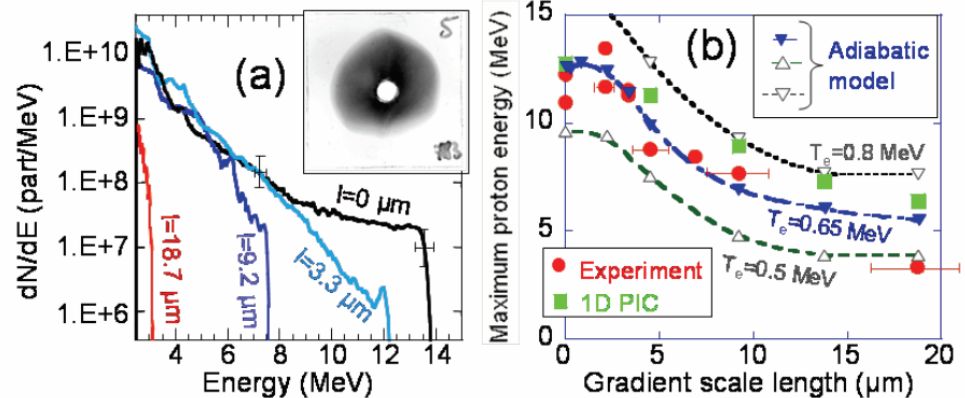
$$E_{\text{max}} = 2 Z T_h \left[\ln(t_p + \sqrt{1+t_p^2}) \right]^2,$$

$$t_p = \omega_{pi} t_{\text{acc}} / \sqrt{2e_N}$$

adiabatic cooling



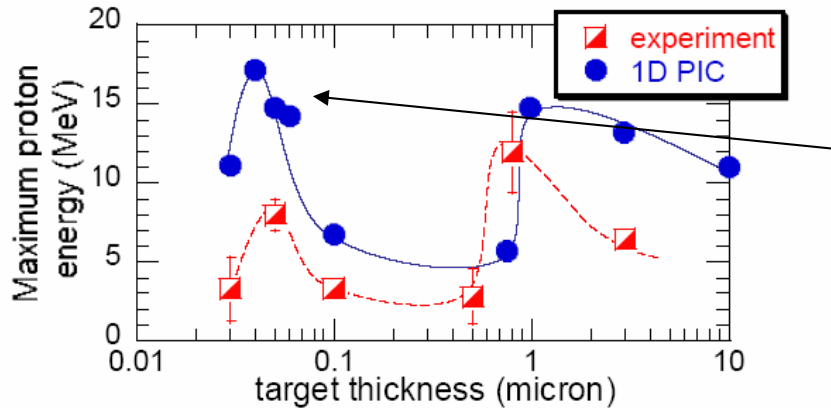
rear surface gradients



Ion acceleration from an underdense plasma

L. Willingale et al., PRL 2006

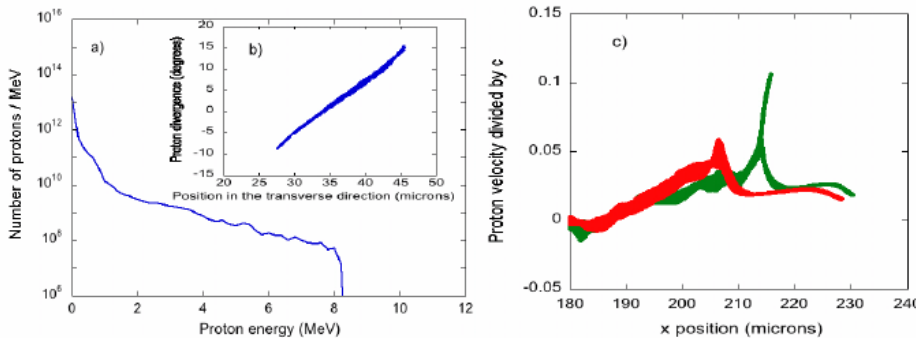
influence of: density gradients (length, shape, amplitude), laser pulse shape



P. Antici et al., NJP 2009

electron heating near the target rear surface

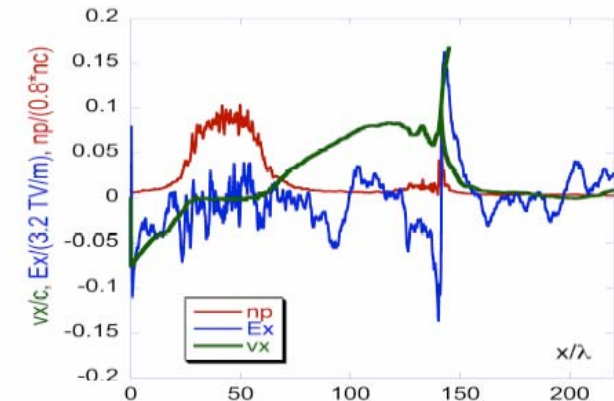
collisionless electrostatic shock



E. d'Humieres, 2010, 2011

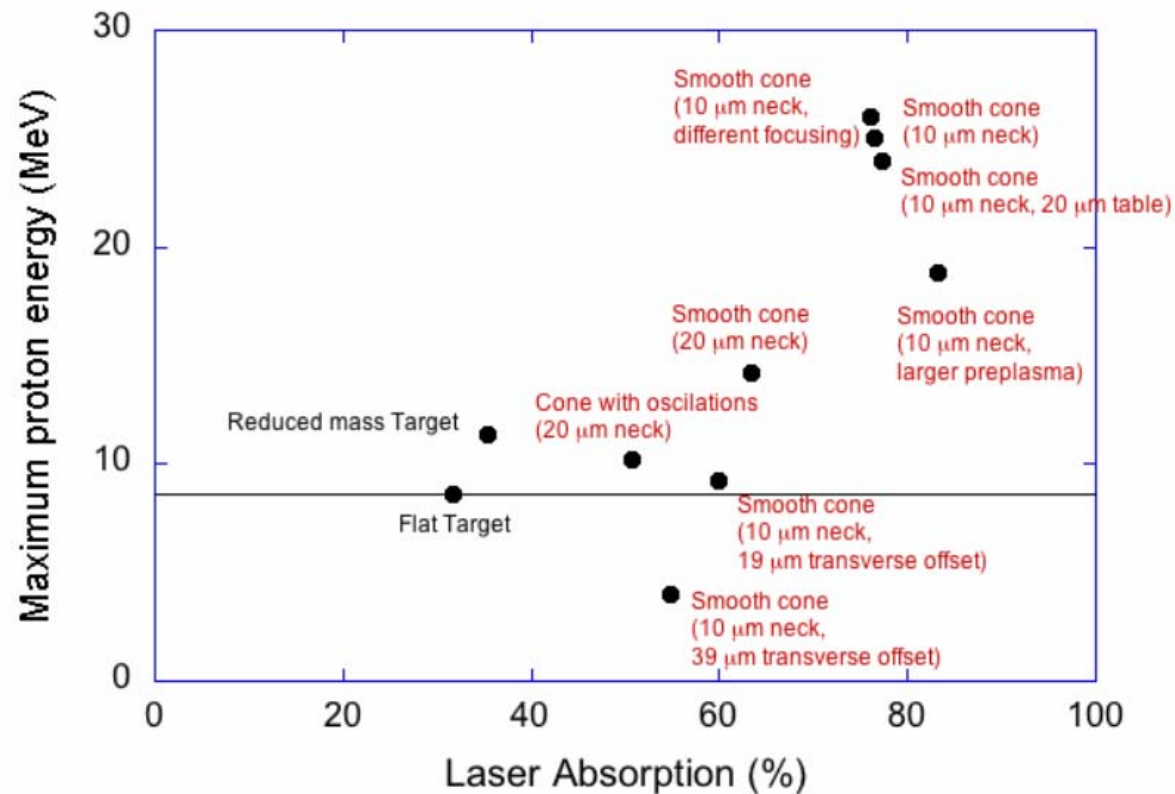
two steps of interaction:

1. electron heating → ion acc.
2. sufficient high ion velocities → reflection of upstreaming ions



Ion acceleration from foils at the tip of a cone

K. Flippo et al., *Phys. Plasmas* 2008

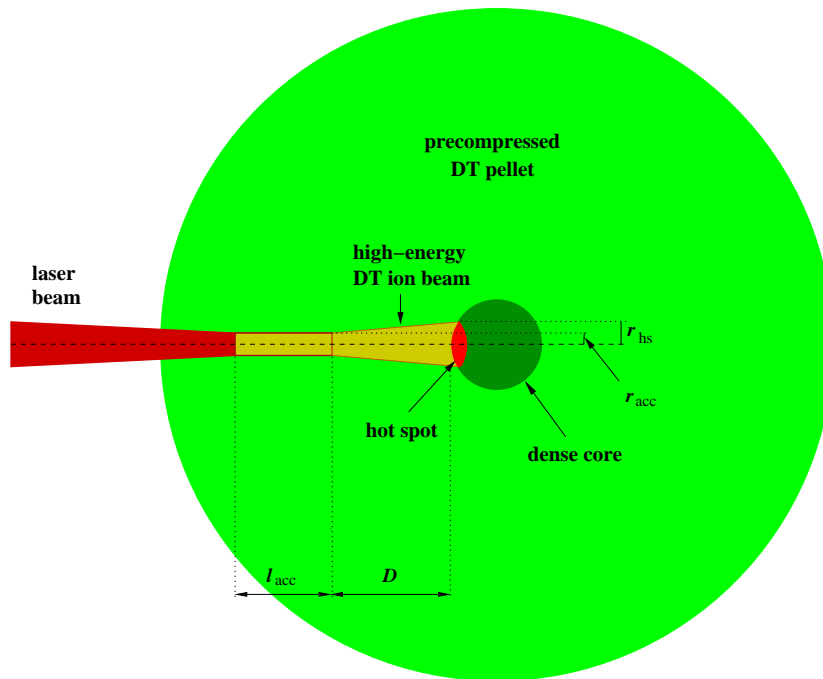


efficient laser absorption, higher T_h and N_h
curved tip surface: cone determines the beam characteristics, divergence control

Alternative fast ion ignition scheme

In-situ DT acceleration

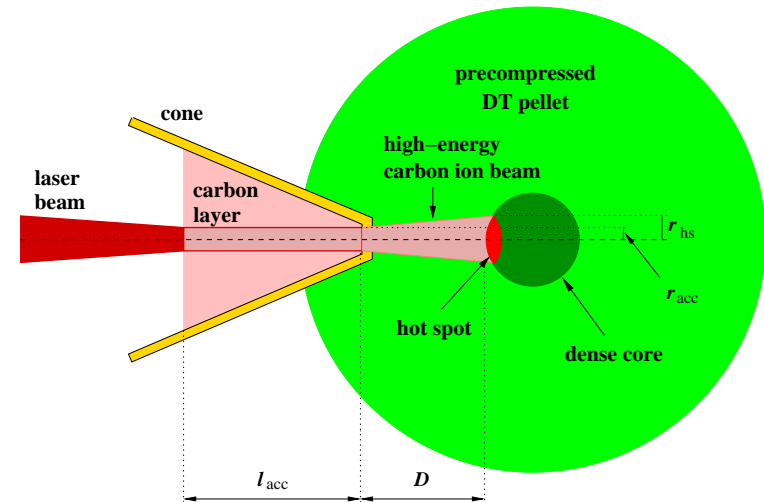
V. Tikhonchuk et al., *Nucl. Fusion* **50**, 045003 (2010)



D.S. Clark and M. Tabak., *Nuclear Fusion* **47**, 1147 (2007)

Cone-guided carbon acceleration

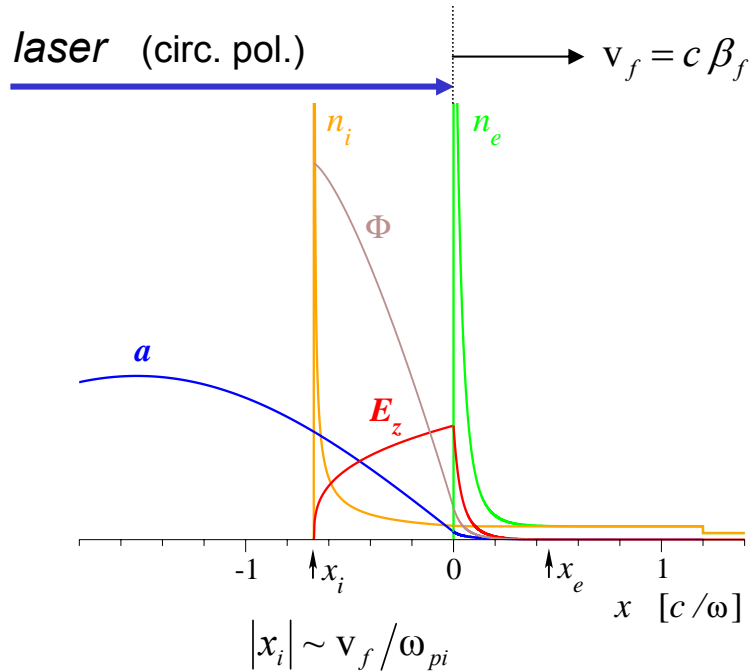
C. Regan et al., *Plasma Phys. Control. Fusion* **53**, 045014 (2011)



A. Debayle et al., *J. Phys. Conf. Ser.* **244**, 022032 (2010)

Analytical piston model

Quasi-stationary laser piston in overdense plasmas



charge separation layer

$$E_z = 2 E_{\text{las}} \sqrt{\frac{1 - \beta_f}{1 + \beta_f}}$$

momentum flux balance:
(in the moving frame MF)

$$\frac{2 I_{\text{las}}}{c} \frac{1 - \beta_f}{1 + \beta_f} = 2 \rho c^2 \gamma_f^2 \beta_f^2$$

$$\rightarrow \beta_f = \frac{B}{1 + B}, \quad B = \sqrt{I_{\text{las}} / \rho c^3}$$

$$\text{LF: } \mathcal{E}_i = 2 m_i c^2 \frac{B^2}{1 + 2B}, \quad 1 - R = \frac{2B}{1 + 2B}$$

main properties confirmed in PIC simulations

A. Robinson et al., *Plasma Phys. Control. Fusion* **51**, 024004 (2009)

N. Naumova et al., *Phys. Rev. Lett.* **102**, 025002 (2009)

T. Schlegel et al., *Phys. Plasmas* **16**, 083103 (2009)

Analytical piston model

Homogeneous plasma

$$\mathcal{E}_i = 2m_i c^2 \frac{B^2}{1+2B}, \quad N_i = \frac{\rho}{m_i} l_{\text{acc}}, \quad F_{\text{acc}} = N_i \mathcal{E}_i, \quad \tau_{\text{las}} = \frac{l_{\text{acc}}}{Bc}, \quad A_{\text{acc}} = \frac{2B}{1+2B}$$

account for beam divergence:

$$F_{\text{hs}} = \eta \frac{N_i \mathcal{E}_i}{\pi r_{\text{acc}}^2 l_{\text{acc}}} \int_D^{D+l_{\text{acc}}} \frac{d\xi}{(1 + \delta\theta \xi / r_{\text{acc}})^2} = \eta F_{\text{acc}} \left[\left(1 + \delta\theta \frac{D}{r_{\text{acc}}}\right) \left(1 + \delta\theta \frac{D+l_{\text{acc}}}{r_{\text{acc}}}\right) \right]^{-1}$$

Plasma density profile with a scale length L

$$N_i = \frac{L}{m_i} (\rho_{\text{max}} - \rho_{\text{min}}), \quad F_{\text{acc}} = \frac{4LI_{\text{las}}}{c} \ln \left(\frac{B_{\text{min}}}{B_{\text{max}}} \frac{1+2B_{\text{max}}}{1+2B_{\text{min}}} \right),$$

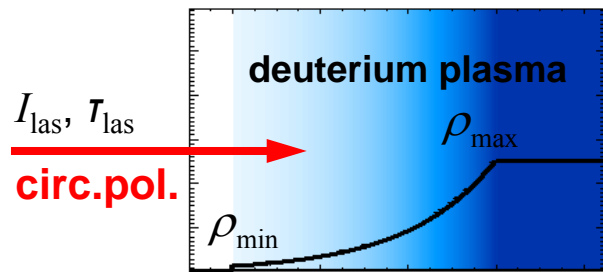
$$\langle \mathcal{E}_i \rangle = F_{\text{acc}} / N_i, \quad \tau_{\text{las}} = \frac{2L}{c} \left(\frac{1}{B_{\text{max}}} - \frac{1}{B_{\text{min}}} \right), \quad A_{\text{acc}} = \frac{F_{\text{acc}}}{I_{\text{las}} \tau_{\text{las}}}$$

1D PIC simulations

Simulation parameters:

$$I_{\text{las}} = 4 \times 10^{22} \text{ W/cm}^2 \quad (a_0 = 100)$$

$$\lambda = 0.8 \text{ } \mu\text{m}$$

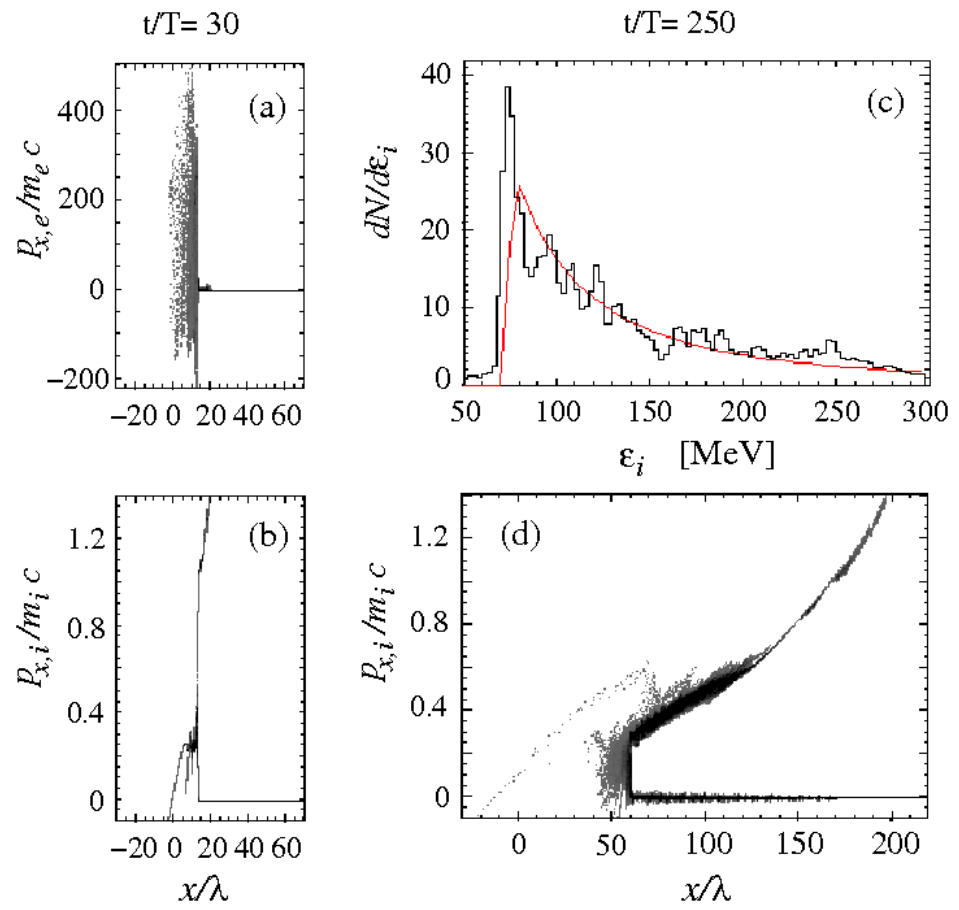


$$L_p = 60 \lambda$$

$$L = \frac{\rho}{d\rho/dx} = 20 \lambda$$

efficiency: $F_{\varepsilon_i=(50-300)\text{MeV}} = 5.4 \text{ GJ/cm}^2$

$$A = F_i / F_{\text{las}} = 0.27$$



2D PIC simulations

Simulation parameters:

$$I_{\text{las}} = 4 \times 10^{22} \text{ W/cm}^2 \quad (a_0 = 100)$$

$$\lambda = 0.8 \mu\text{m}$$

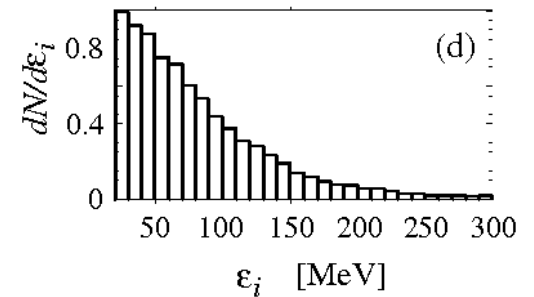
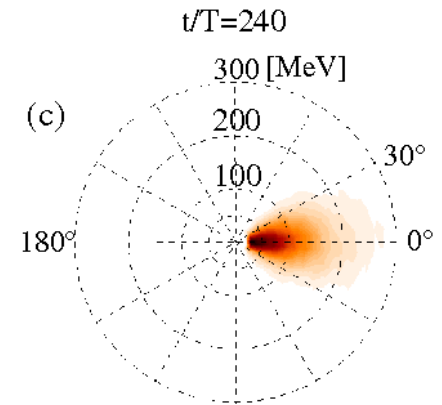
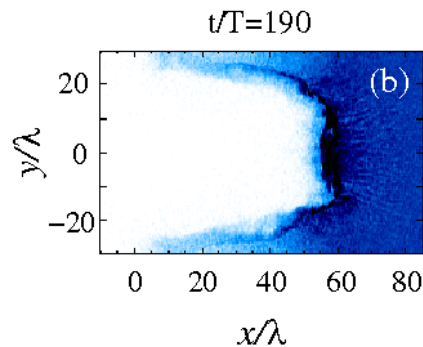
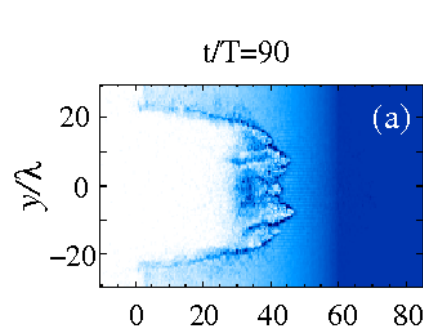
flat-top transverse laser-intensity profile: $w = 20 \lambda$

deuterium plasma

clean and stable channel

small angular spread of ions

$$\varepsilon_i \geq 50 \text{ MeV}: \quad \theta \leq 5^\circ$$



Requirements for fast ignition application

- Circular laser polarization
- Overcritical plasma: $n_e / a_0 n_c \geq 1$
- Large number of accelerated ions, $\mathcal{N}_i \sim 10^{14} \div 10^{16}$, to provide: $F_{\text{hs:ig}} \geq E_{\text{ig}} / \pi r_{\text{ig}}^2$
- Sufficiently small laser focus, $r_{\text{acc}} < r_{\text{ig}}$, because of ion beam divergence
- Rayleigh length of the laser pulse $>$ acceleration length l_{acc} } for mono-energetic ion acceleration
- High laser contrast, flat transverse intensity profile

IFI scheme design

1. Essential precompressed pellet parameters: ρ_{\max} , $(\rho R)_{\text{core}}$, L

2. Ignition scaling laws $\rightarrow E_{\text{ig}}, r_{\text{ig}}, \tau_{\text{ig}}$ S. Atzeni, *Phys. Plasmas* **6**, 3316 (1999)

③ Supposed density ratio of the acceleration zone, $\rho_{\max} / \rho_{\min}$, and D
 $\rightarrow l_{\text{acc}}, \rho_{\max}, \rho_{\min}, N_i$

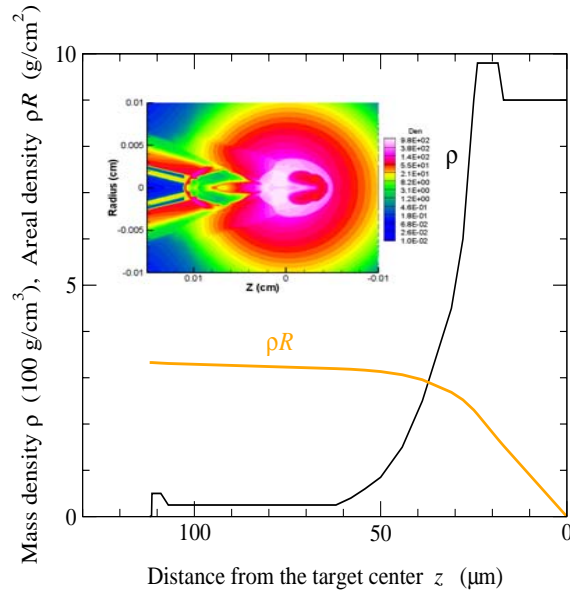
④ With a chosen laser intensity $I_{\text{las}} \rightarrow F_{\text{acc}}, \langle \mathcal{E}_i \rangle$

5. Laser focal spot radius r_{acc} and ion beam divergence $\delta\theta$ determine F_{hs}

$$\eta F_{\text{hs:ig}} \pi r_{\text{ig}}^2 \propto E_{\text{ig}}$$

6. Final estimate of $P_{\text{las}}, \tau_{\text{las}}, E_{\text{las}}, A_{\text{acc}}$

Cone-guided target



Approximate DT density profile: $(\rho R)_{\text{core}} \approx 3 \text{ g/cm}^2$

Carbon layer with $\rho = 0.2 \text{ g/cm}^3$

Scalings for ignition:

$$E_{\text{ig}} = 7 \text{ kJ}, \quad r_{\text{ig}} = 12 \mu\text{m}, \quad \tau_{\text{ig}} = 14 \text{ ps}$$

$$\longrightarrow F_{\text{ig}} \approx 1.6 \text{ GJ/cm}^2$$

Estimate:

$$\varepsilon_i = 450 \text{ MeV} \quad \longrightarrow \quad I_{\text{las}} = 1.4 \times 10^{22} \text{ W/cm}^2 \quad (a_0 \approx 60)$$

$$r_{\text{acc}} = 6 \mu\text{m}, \quad l_{\text{acc}} = 200 \mu\text{m}$$

$$\longrightarrow N_i = 2 \times 10^{20} \text{ cm}^{-2}, \quad F_{\text{acc}} \approx 14 \text{ GJ/cm}^2$$

$$D = 60 \mu\text{m}, \quad \delta\theta = 0.1: \quad r_{\text{hs}}^2 \approx 11 r_{\text{acc}}^2, \quad F_{\text{hs}} \approx 1.3 \text{ GJ/cm}^2$$

$$E_{\text{hs:ig}} = 5 \text{ kJ} \quad (\eta = 0.8) \quad \longrightarrow \quad \text{no ignition}$$

$$N_i = \frac{\rho}{m_i} l_{\text{acc}}, \quad F_{\text{acc}} = N_i \varepsilon_i,$$

$$\tau_{\text{las}} = \frac{l_{\text{acc}}}{Bc}, \quad A_{\text{acc}} = \frac{F_{\text{acc}}}{I_{\text{las}} \tau_{\text{las}}}$$

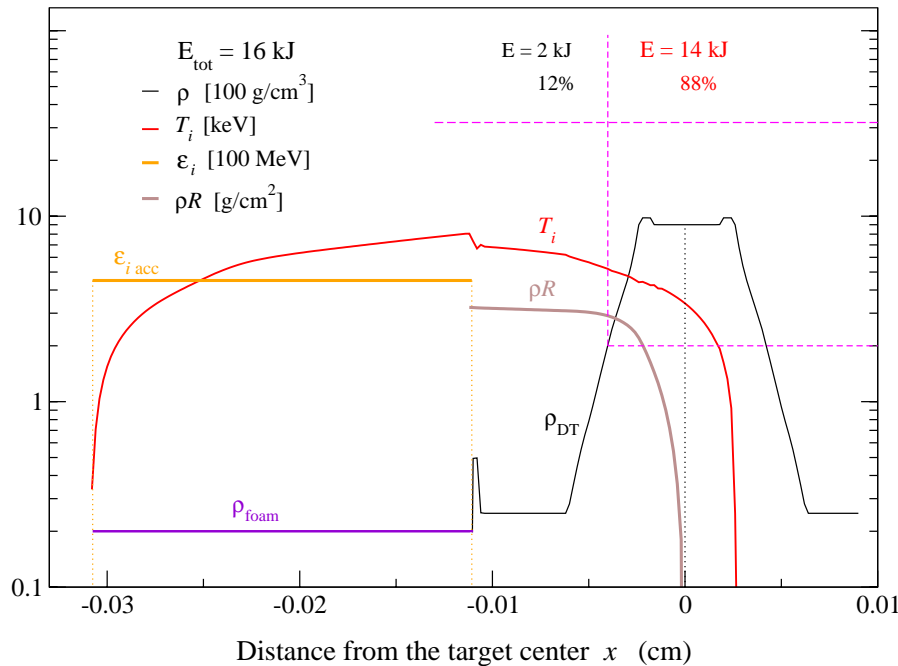
Higher directionality

$$\delta\theta = 0.05$$

$$\longrightarrow F_{\text{hs:ig}} = 3 \text{ GJ/cm}^2$$

$$E_{\text{hs:ig}} = 11 \text{ kJ } (\eta = 0.8)$$

Profiles after ion stopping:
(*Trumpet code*)



successful ignition

$$P_{\text{las}} = 16 \text{ PW}, \tau_{\text{las}} = 4 \text{ ps}$$

$$E_{\text{las}} = 65 \text{ kJ}$$

$$A_{\text{acc}} \approx 25\%$$

$$A_{\text{ig}} = \frac{\eta F_{\text{hs}} \pi r_{\text{ig}}^2}{P_{\text{las}} \tau_{\text{las}}} \approx 17\% (\eta = 0.8)$$

Reduced laser intensity

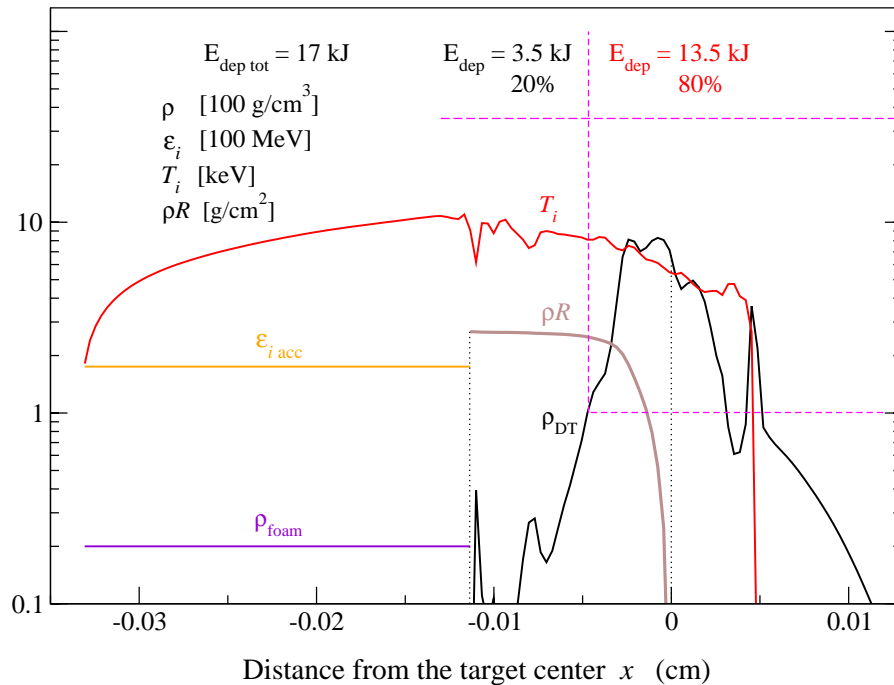
Scenario with lower laser intensity

$$I_{\text{las}} = 5 \times 10^{21} \text{ W/cm}^2 \rightarrow \mathcal{E}_i = 175 \text{ MeV}$$

$$(a_0 \approx 35)$$

Profiles after ion stopping:
(Trumpet code)

Design with same D , l_{acc} , $\delta\theta$ as before:



With $r_{\text{acc}} \approx 10 \mu\text{m}$:

$$F_{\text{acc}} = 5.6 \text{ GJ/cm}^2$$

$$F_{\text{hs}} = 1.9 \text{ GJ/cm}^2$$

$$E_{\text{hs:ig}} = 8.4 \text{ kJ}$$

ignition

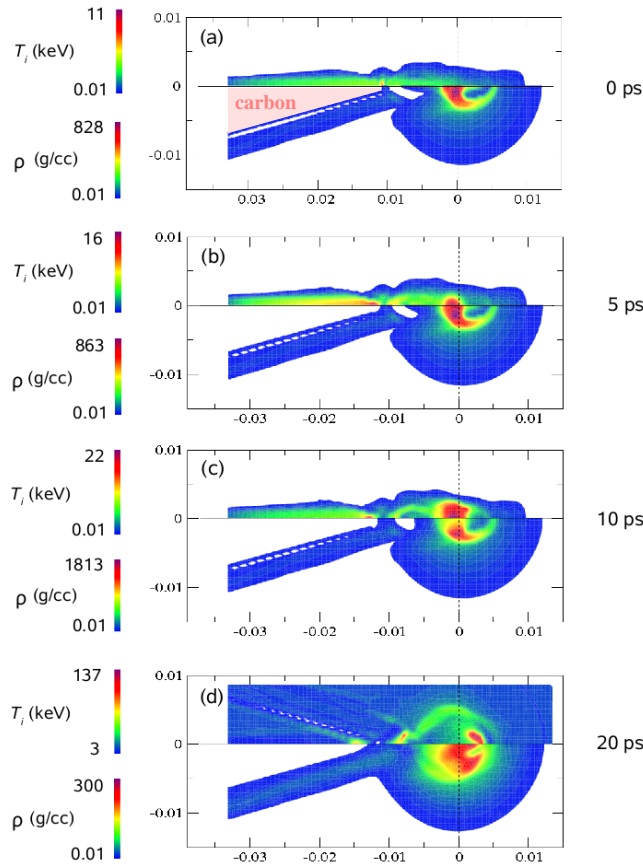
$$P_{\text{las}} = 15 \text{ PW}, \tau_{\text{las}} = 7 \text{ ps}$$

$$E_{\text{las}} = 100 \text{ kJ}$$

$$A_{\text{acc}} \approx 17\%$$

DT target dynamics

2D axially-symmetric CHIC simulation after ion beam energy deposition

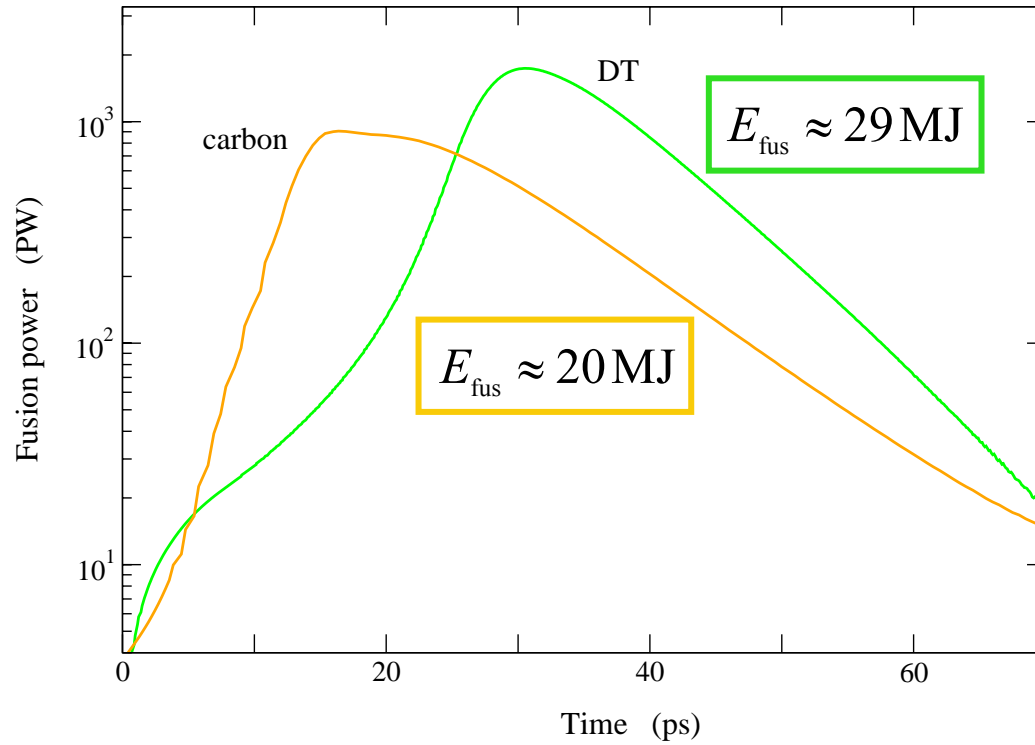


Further heating after ion beam energy deposition by α - particles

$$\text{reaction rate: } \nu_{\text{DT}} \approx 2.6 \times 10^5 \rho T_i^2 \text{ s}^{-1}$$

$$\text{equilibration time: } \tau_{\alpha e} \approx \beta \frac{T_e^{3/2}}{\rho} \text{ ps}$$

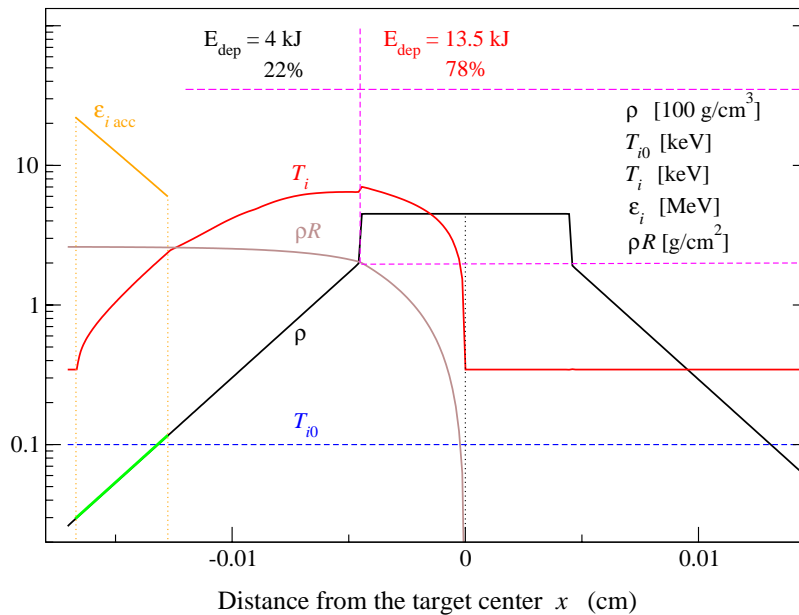
Fusion power and energy



In-situ DT acceleration

Large-scale target: $(\rho R)_{\text{core}} \approx 2.6 \text{ g/cm}^2$, $\rho_{\text{max}} \approx 450 \text{ g/cm}^3$

Profiles after ion stopping:
(*Trumpet code*)



Ignition scaling laws:

$$E_{\text{ig}} = 8.5 \text{ kJ}, \quad r_{\text{ig}} = 13.5 \mu\text{m}, \quad \tau_{\text{ig}} = 15 \text{ ps}$$

DT acceleration: $\langle \epsilon_i \rangle \approx 10 \text{ MeV}$

$$I_{\text{las}} = 4.6 \times 10^{22} \text{ W/cm}^2, \quad r_{\text{acc}} = 7 \mu\text{m}$$

$$l_{\text{acc}} = 40 \mu\text{m}, \quad D = 80 \mu\text{m}$$

$$N_i \approx 6 \times 10^{21} \text{ cm}^{-2}, \quad F_{\text{acc}} \approx 11 \text{ GJ/cm}^2$$

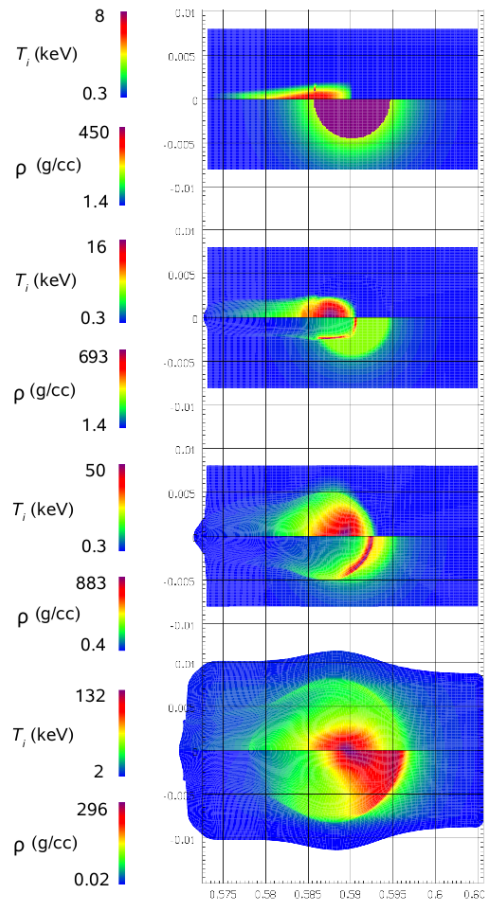
Beam divergence: $\delta\theta = 0.1$

$$\rightarrow F_{\text{hs:ig}} = 1.9 \text{ GJ/cm}^2$$

$$E_{\text{hs:ig}} = 9 \text{ kJ} \quad (\eta = 0.2)$$

Ignition requirements: $P_{\text{las}} = 70 \text{ PW}$, $\tau_{\text{las}} = 2.5 \text{ ps}$, $E_{\text{las}} = 180 \text{ kJ}$, $A_{\text{acc}} \approx 9\%$

DT target dynamics



2D axially-symmetric CHIC simulation
after ion beam energy deposition

0 ps

10 ps

26 ps

36 ps

← ignition

← combustion wave

Conclusions

High-energy carbon ions

- high acceleration efficiency ← *lower plasma density*
- enhanced localization of ion energy deposition
- low ion beam divergence ← *high ion energy*
- reduced demand of laser energy for DT ignition
- ion beam divergence?
- cone

In-situ DT acceleration

- simple geometry
- but: lower acceleration efficiency
 - larger ion beam divergence
 - hole-boring through undercritical plasma
 - higher laser power and energy